

Interplay of magnetic and electronic states in the pyrochlore iridates

Leon Balents, KITP, UCSB



- or -
correlations versus
topology

APS meeting, Baltimore, March 2016

People



Ru Chen



Lucile Savary



Eun-Gook Moon



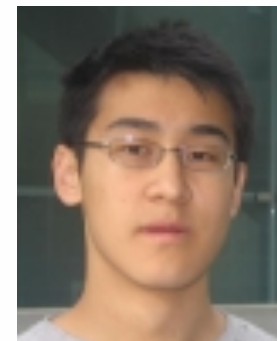
S. Nakatsuji



T. Kondo



Z. Tian



T. Hsieh



H. Ishizuka

General questions about pyrochlore iridates

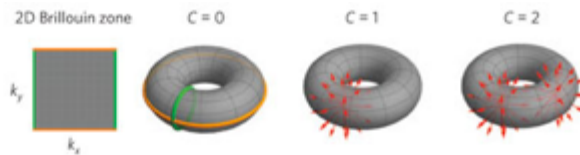
- How correlated are these materials?
 - Lots of theory uses mean-field. Does this work? If so why?
- What is the role of interactions beyond mean-field? Slater versus Mott.
- What is the nature of the MIT?
- Any real signatures of topology?

TIs

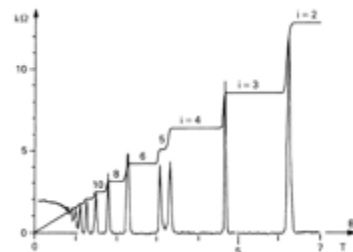
Broken
TRS

with
TRS

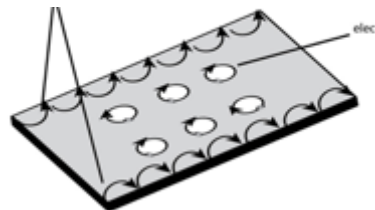
integer Chern
number



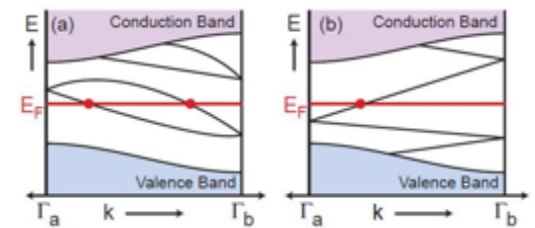
= Hall
conductance



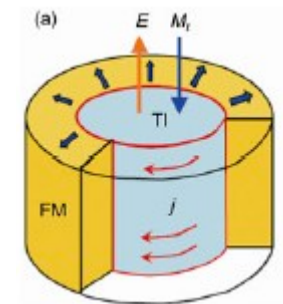
chiral edge
states



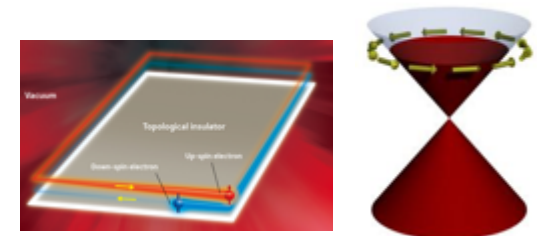
Z_2
invariant



= M-E
polarizability



helical edge/
surface states

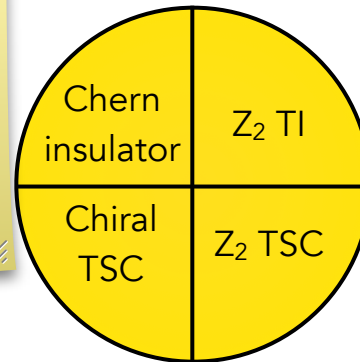
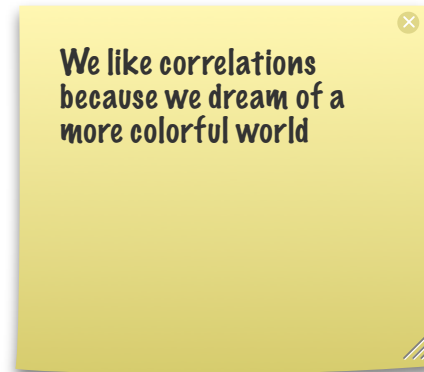


Topological everything

- Chern insulator
- Z_2 topological insulator
- topological crystalline insulators
- topological superconductors
- Dirac semi-metals
- Weyl semi-metals
- Line node semi-metals
- Fermi surface
- ...

Why correlations?

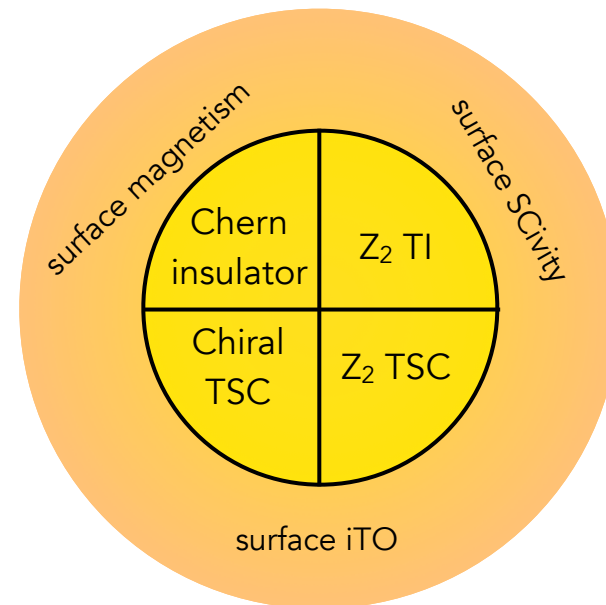
w/o interactions



+ Weyl, Dirac...

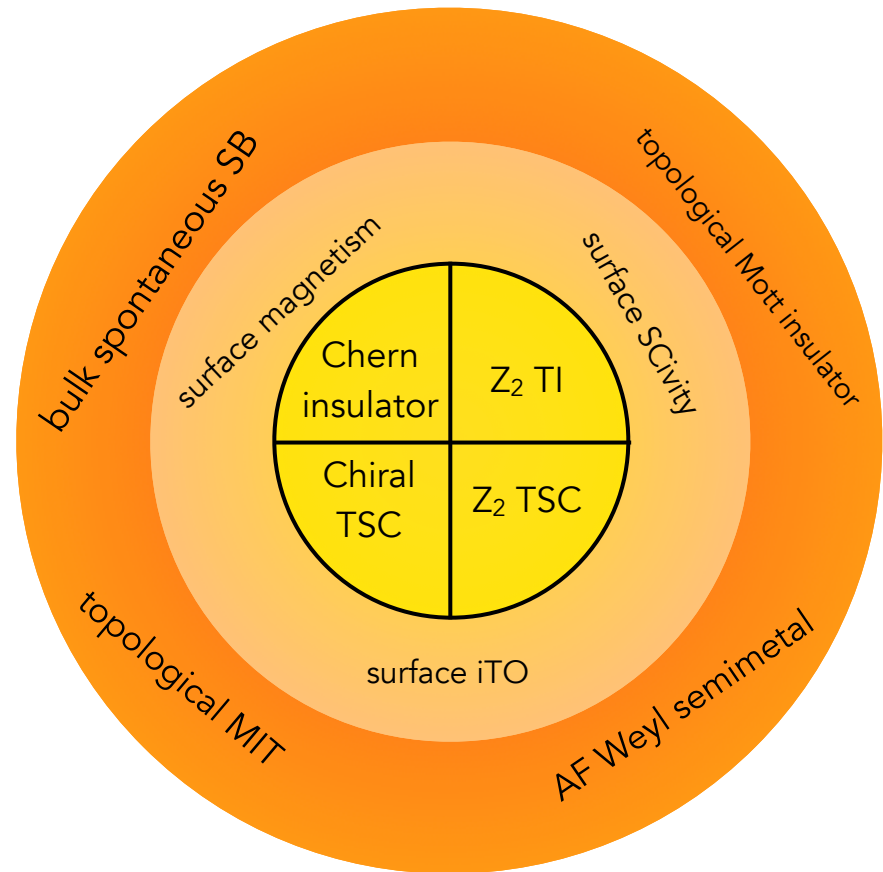
Why correlations?

- Distinct surface states



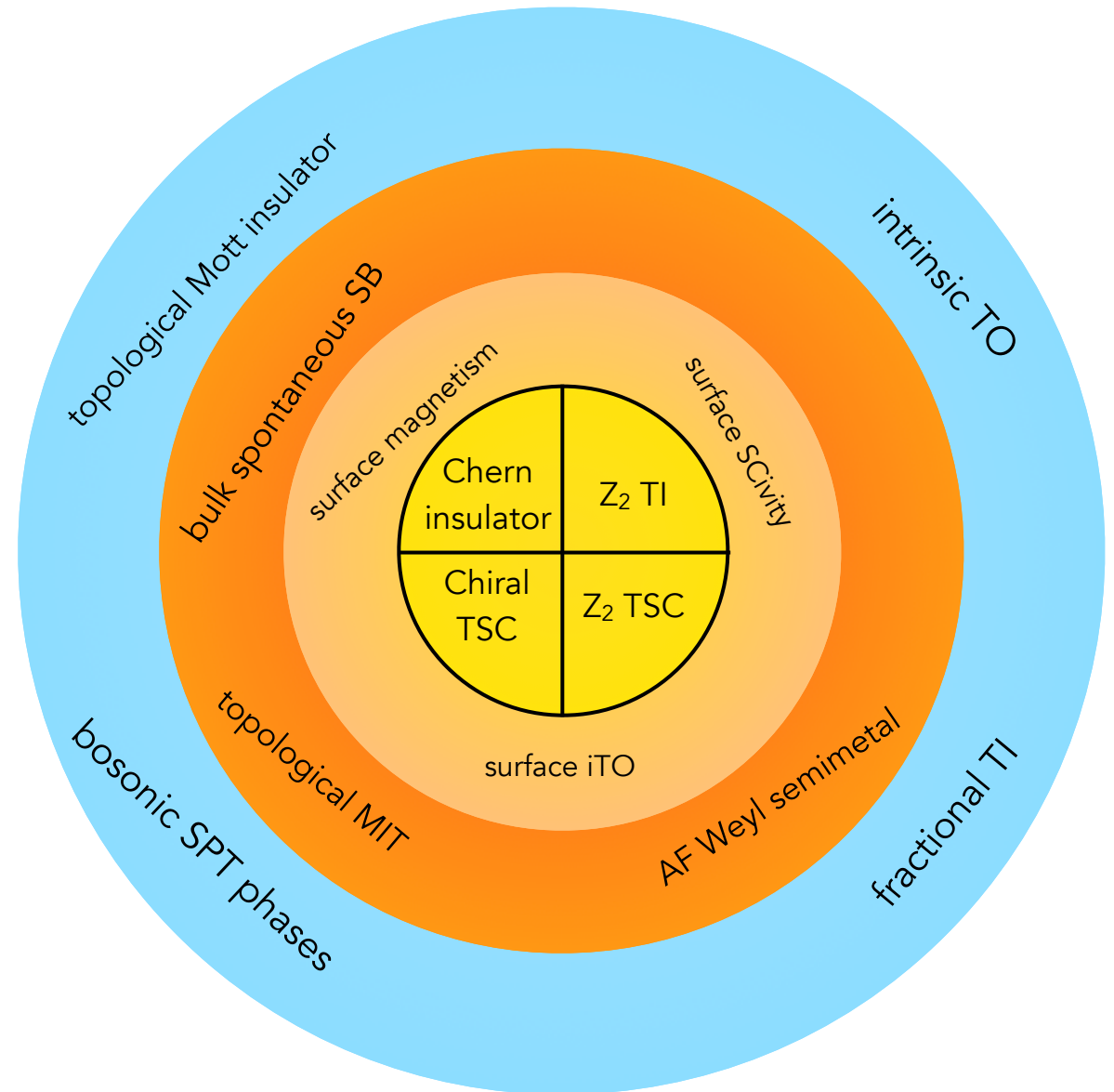
Why correlations?

- Distinct surface states
- Control of bulk topology by spontaneous symmetry breaking



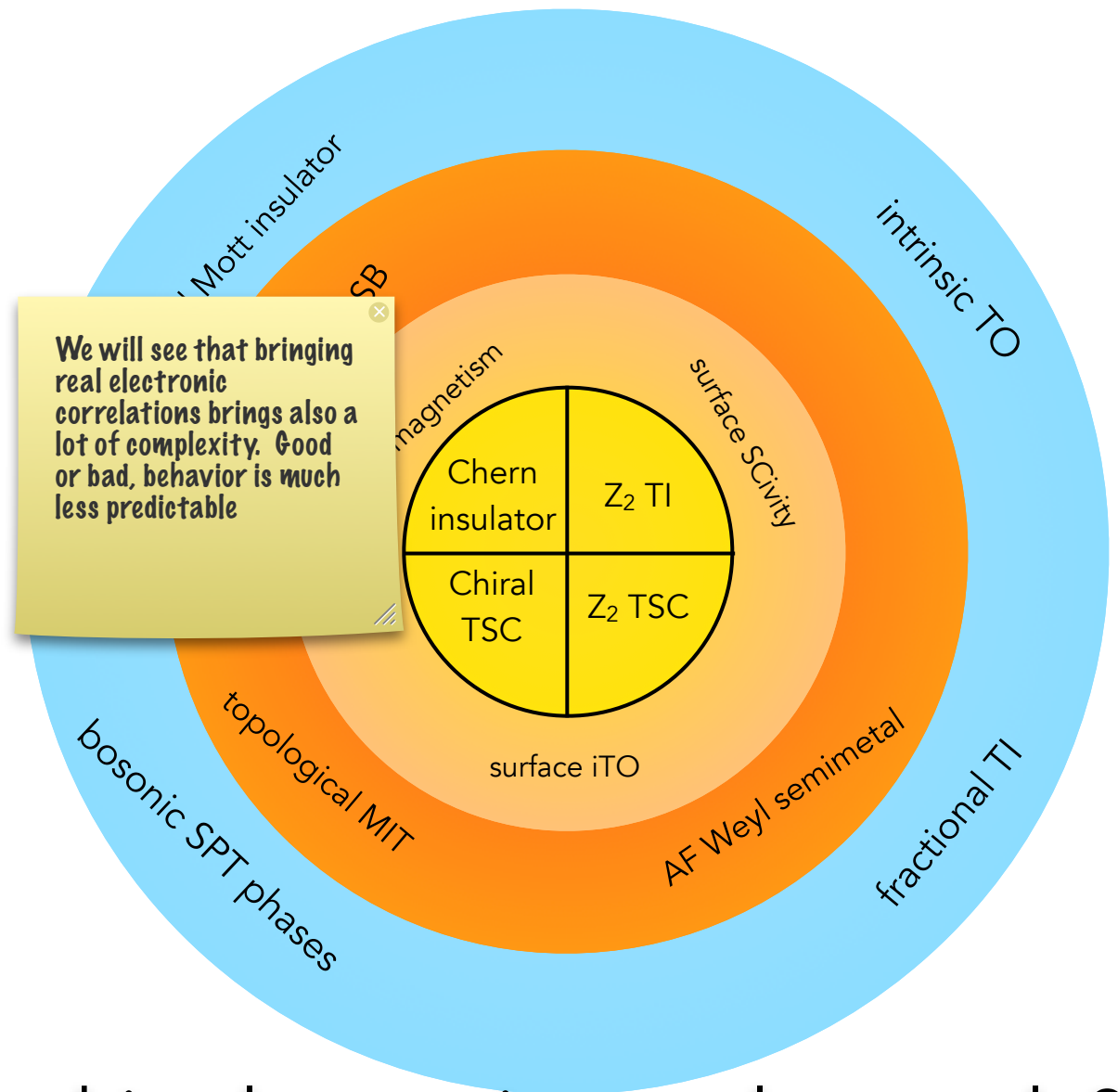
Why correlations?

- Distinct surface states
- Control of bulk topology by spontaneous symmetry breaking
- Entirely new phases requiring strong interactions



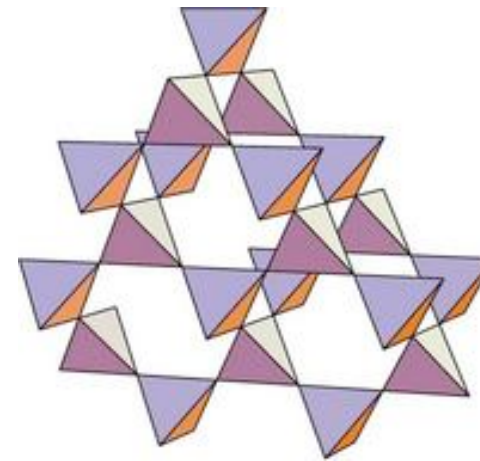
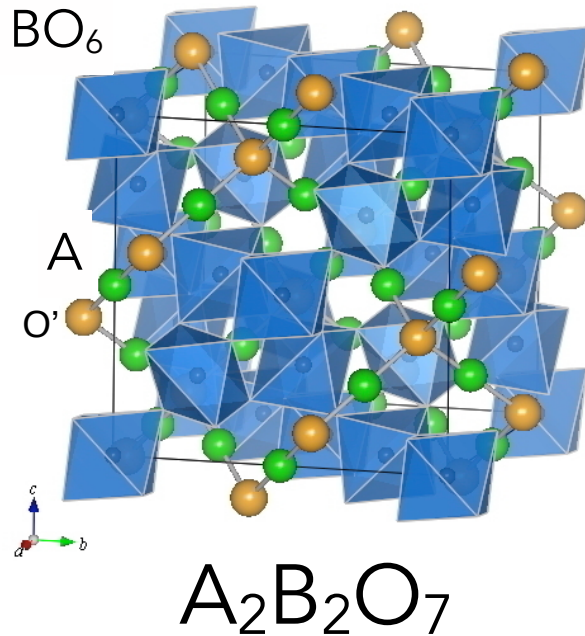
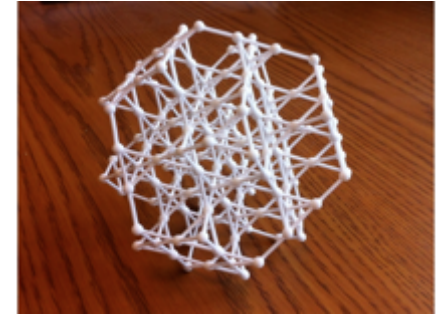
Why correlations?

- Distinct surface states
- Control of bulk topology by spontaneous symmetry breaking
- Entirely new phases requiring strong interactions



This talk: how does this play out in a real example?

Pyrochlores

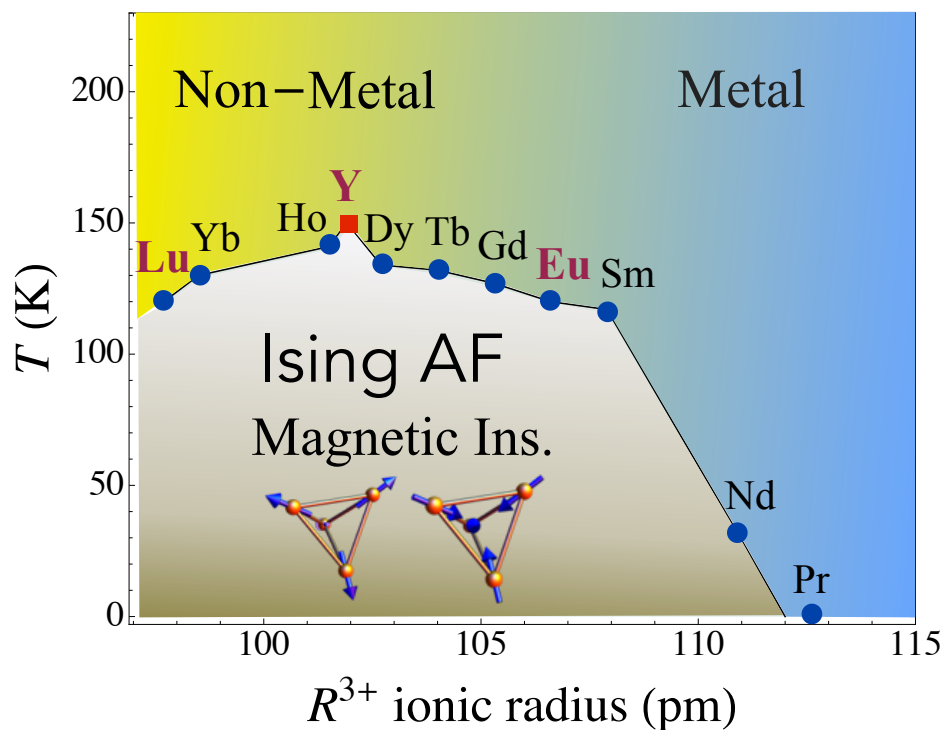


A (B) sublattice

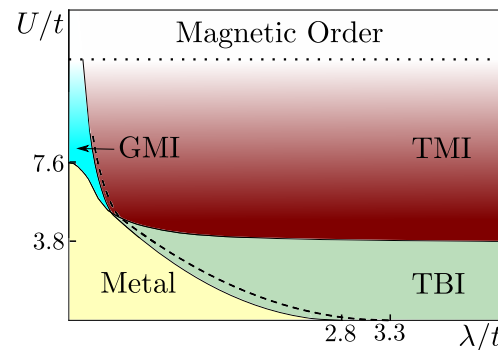
$A = Y, Ln$ ← f-electron moments, except
 $B = Ir$ when $A = Y, Eu$

Pyrochlore iridates

- Continuous magnetic/metal-insulator transitions (compatible with Ising)



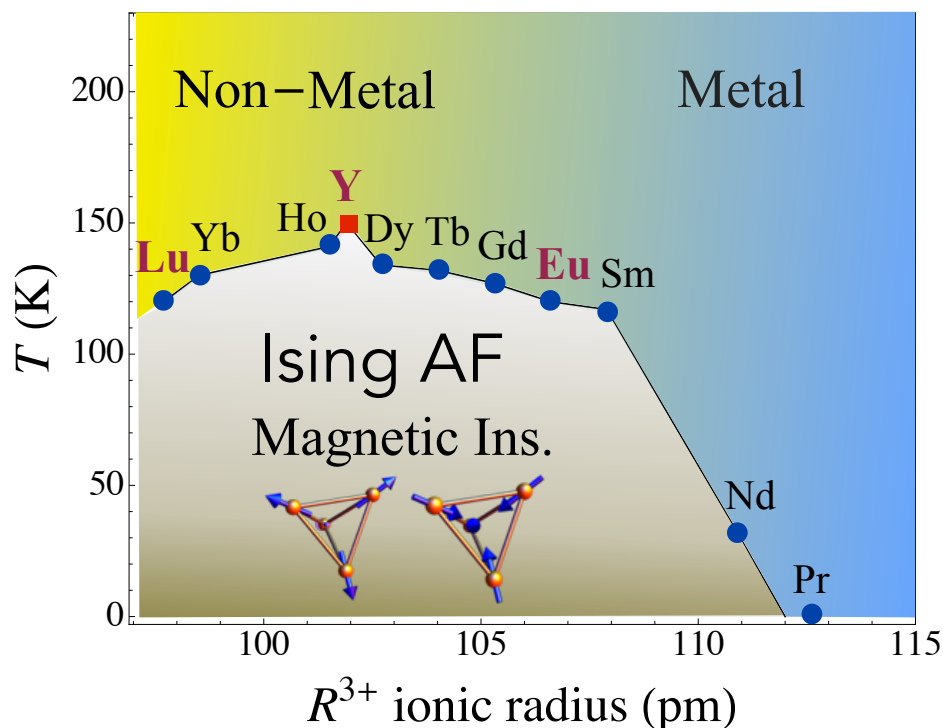
Yanagashima+Maeno, JPSJ 2001
 K. Matsuhira et al, JPSJ 2011
 W. Witczak-Krempa et al, ARCOMP 2013



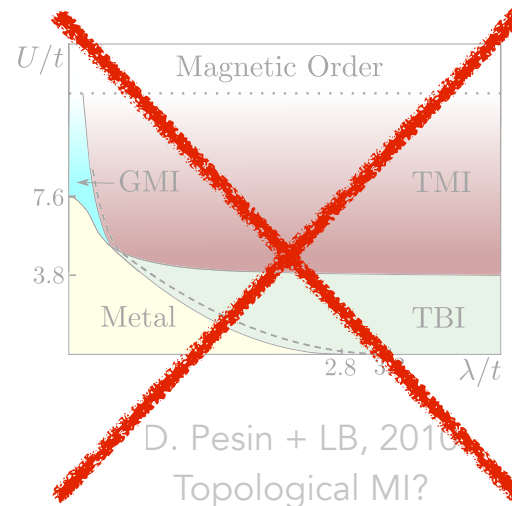
D. Pesin + LB, 2010
 Topological MI?

Pyrochlore iridates

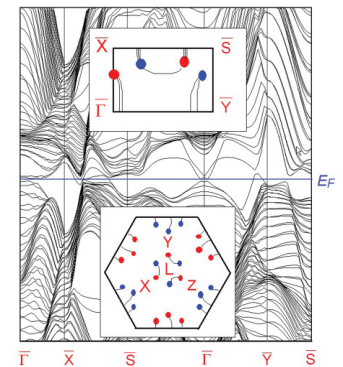
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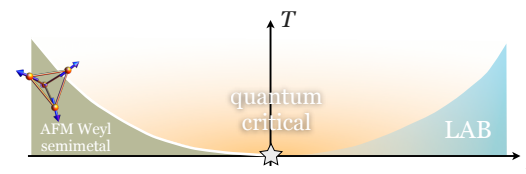
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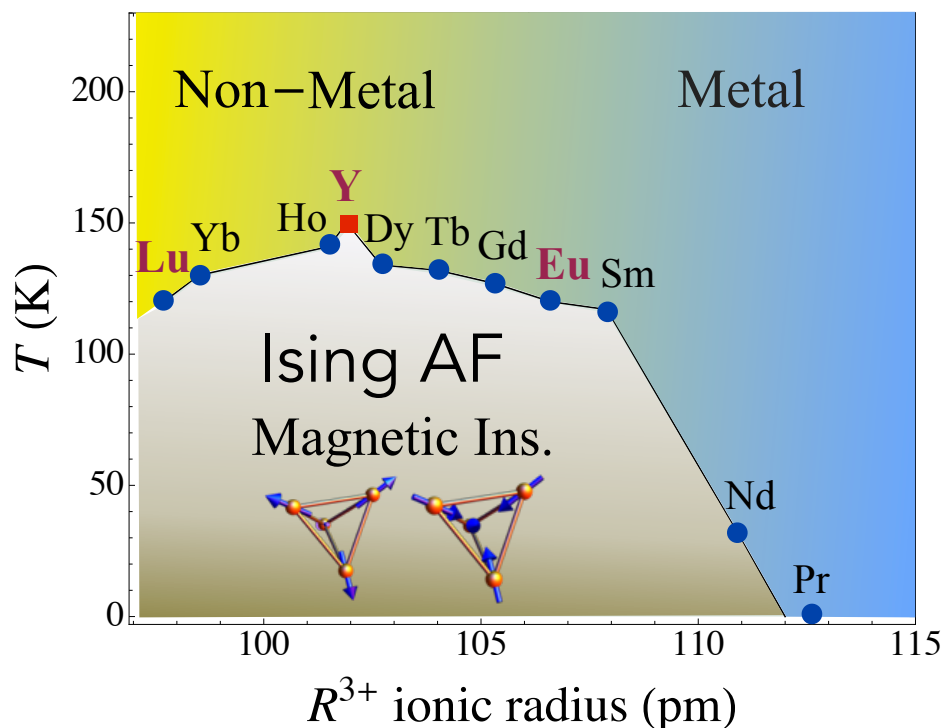
X. Wan et al, 2011
AF Weyl semimetal?



L. Savary et al, 2014 - topological QCP?

Pyrochlore iridates

- Continuous magnetic/metal-insulator transitions (compatible with Ising)



R^{3+} ionic radius (pm)

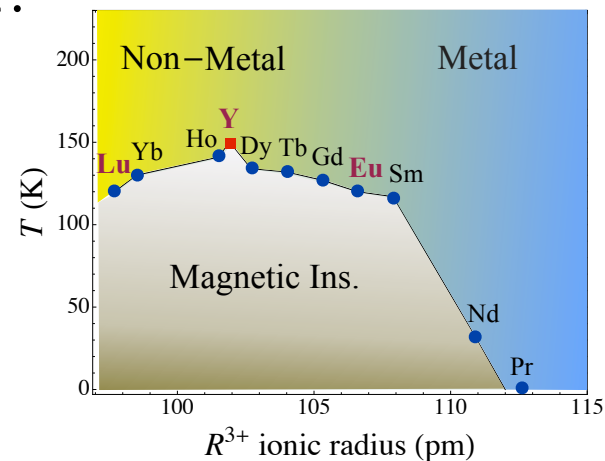
Yanagashima+Maeno, JPSJ 2001
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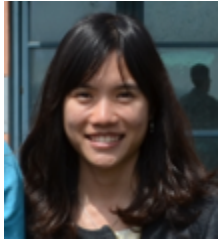
How correlated are these materials? Is there physics beyond mean-field? If yes, what is its role?

Ingredients of a theory

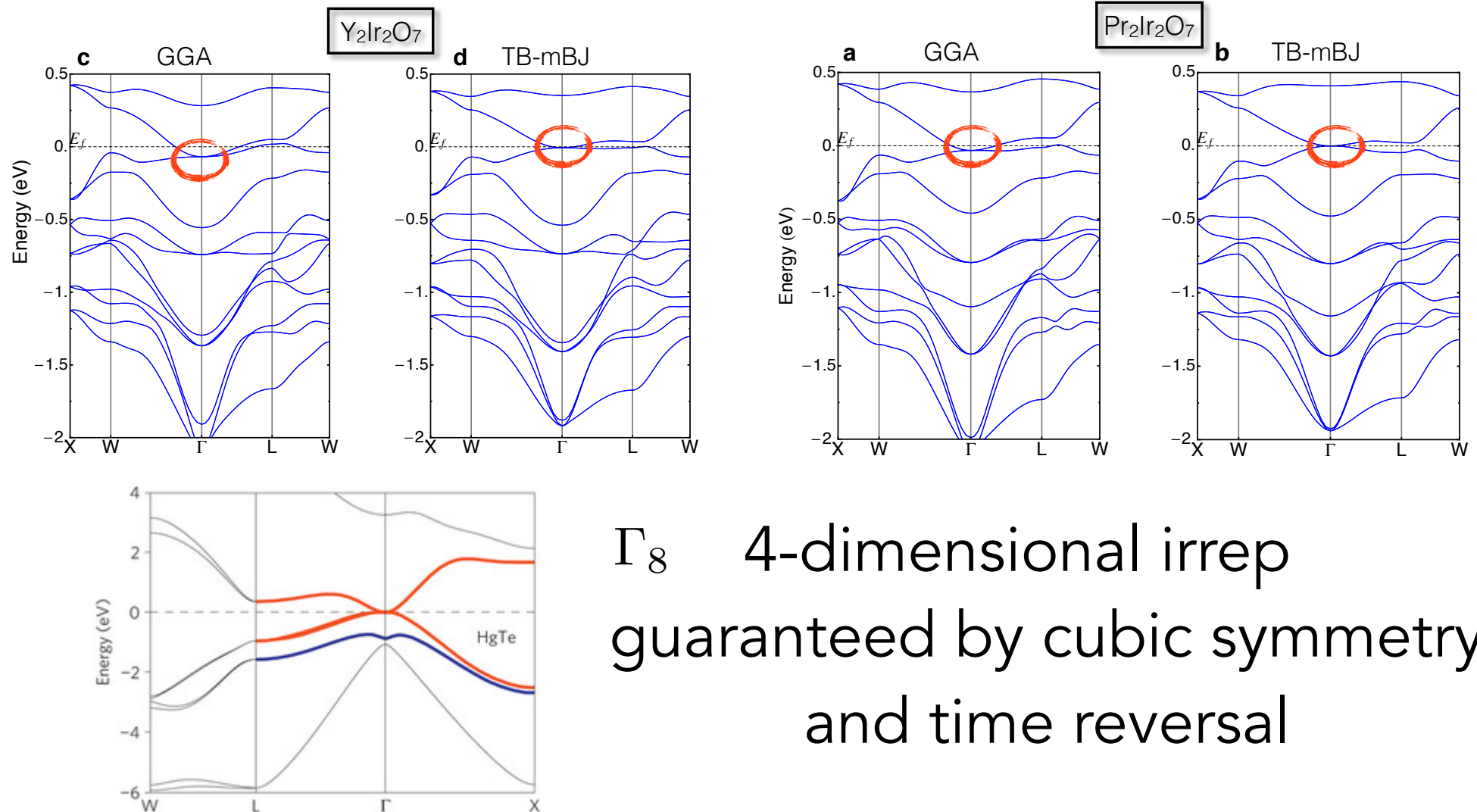
- Unusual electronic structure
- Ir e-e Hubbard interactions
- Rare earth moments?

probably not?

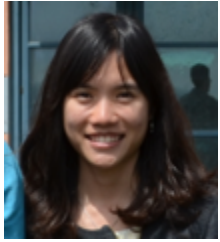




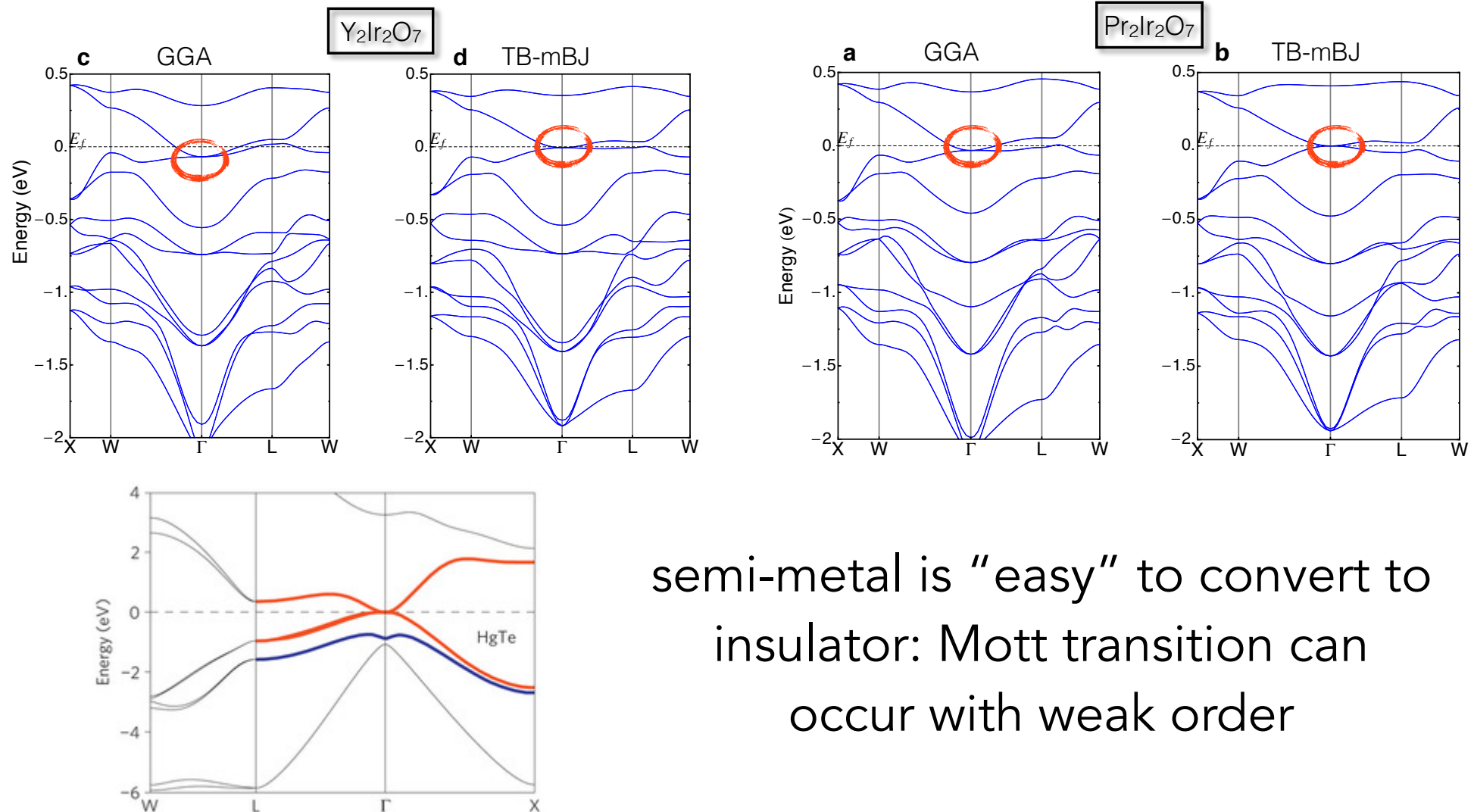
Paramagnetic electronic structure



Γ_8 4-dimensional irrep
guaranteed by cubic symmetry
and time reversal

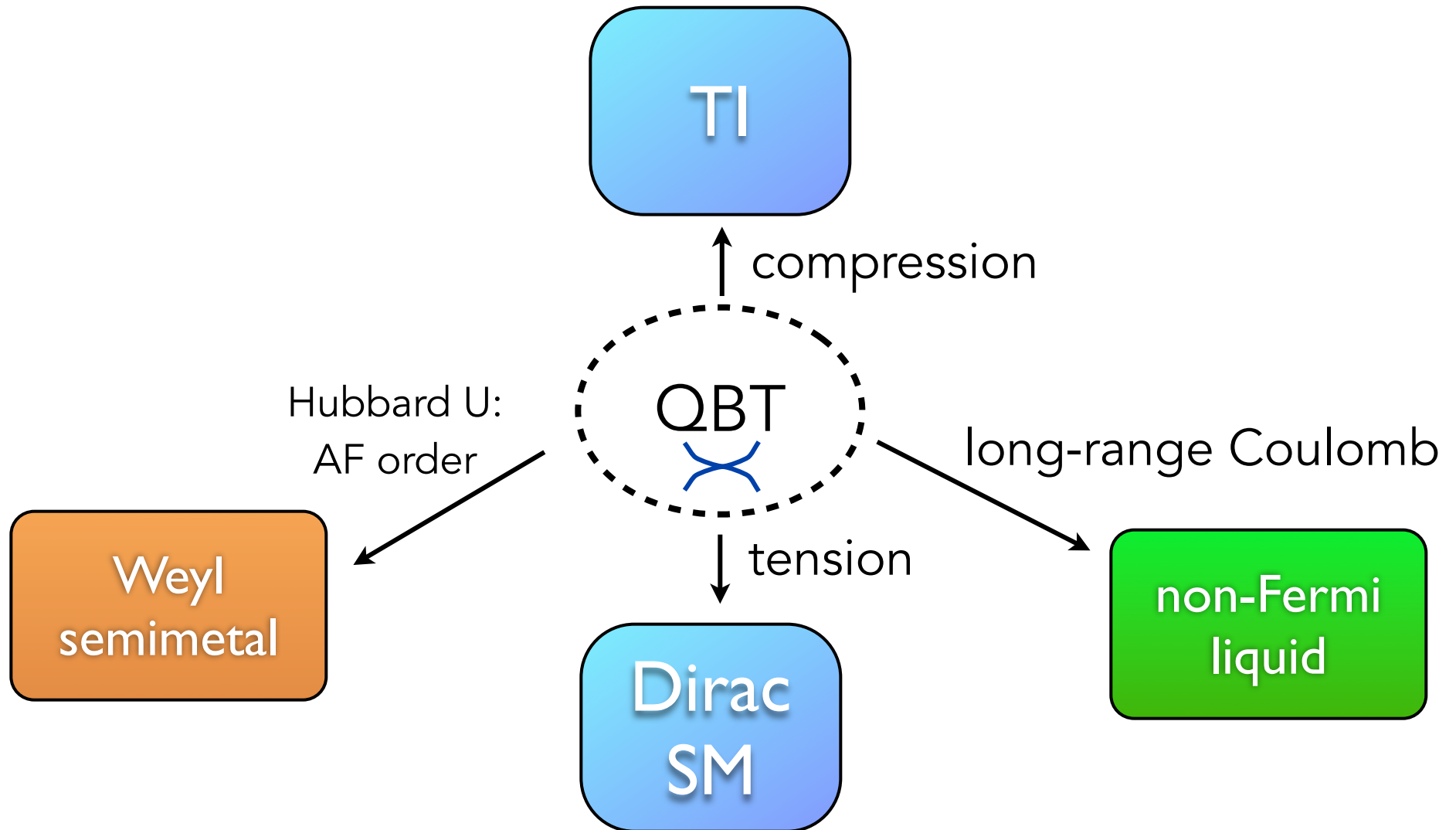


Paramagnetic electronic structure



semi-metal is "easy" to convert to insulator: Mott transition can occur with weak order

Proximate phases





ARPES

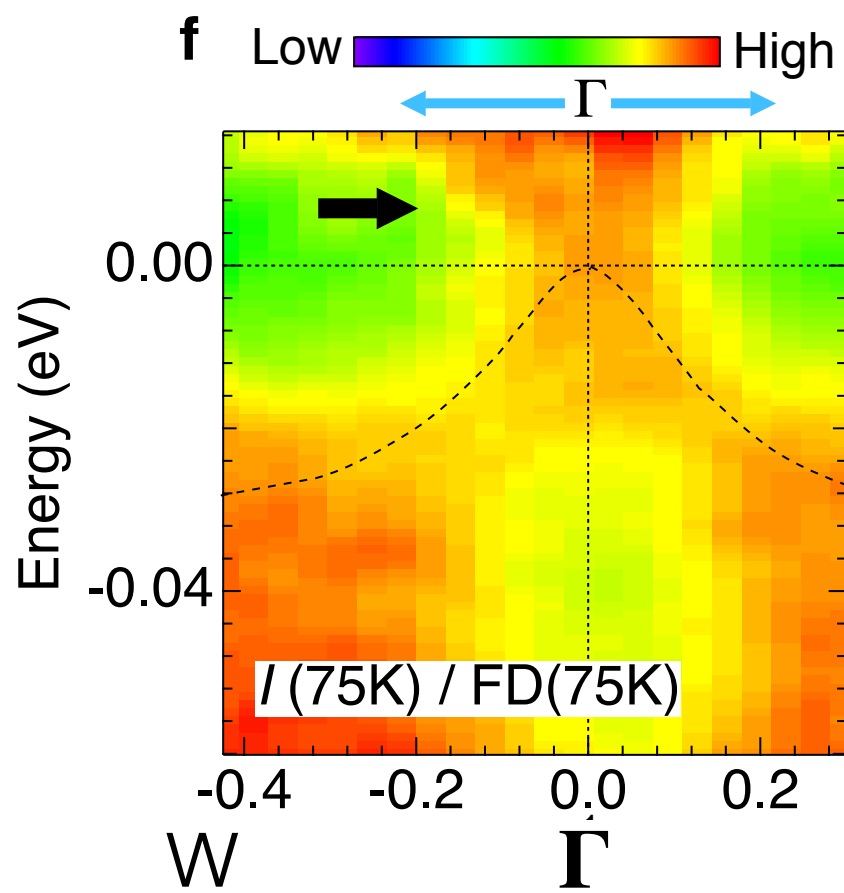
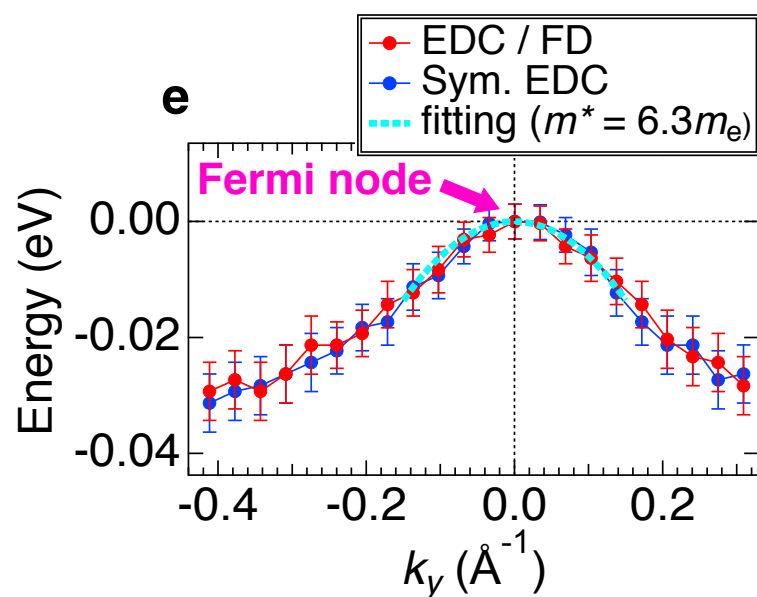
T. Kondo et al, Nat. Comm., 2015



$\text{Pr}_2\text{Ir}_2\text{O}_7$ S. Nakatsuji

T. Kondo

S. Shin





ARPES

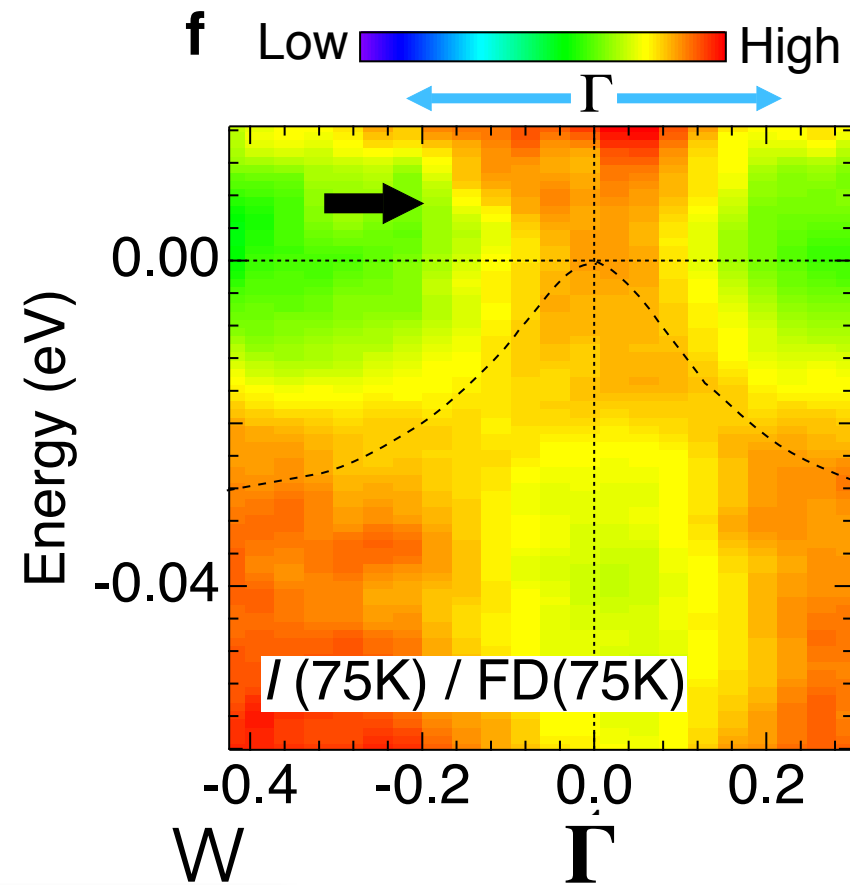
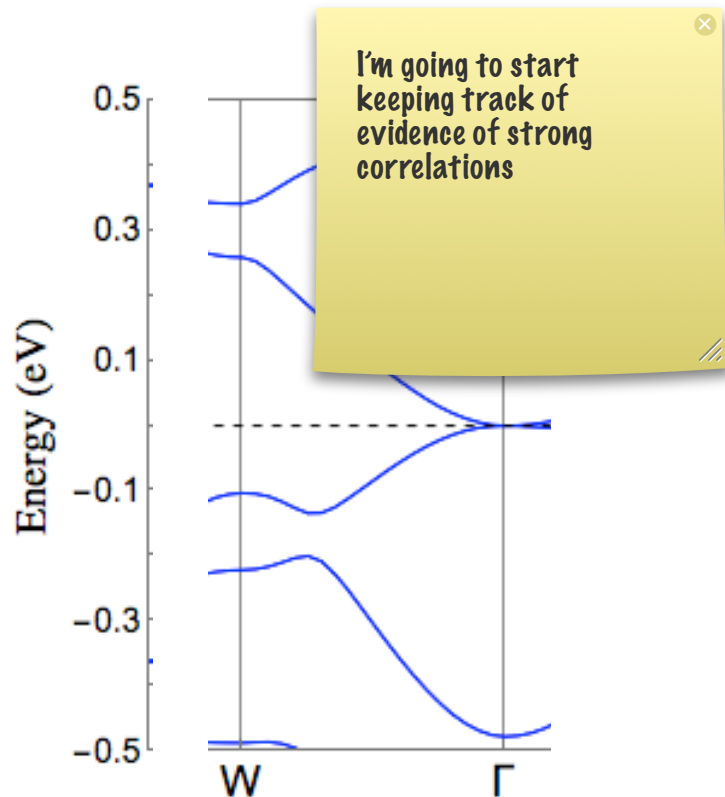
T. Kondo et al, Nat. Comm., 2015



$\text{Pr}_2\text{Ir}_2\text{O}_7$ S. Nakatsuji

T. Kondo

S. Shin



✓ Bandwidth reduced by 3-5 from DFT



ARPES

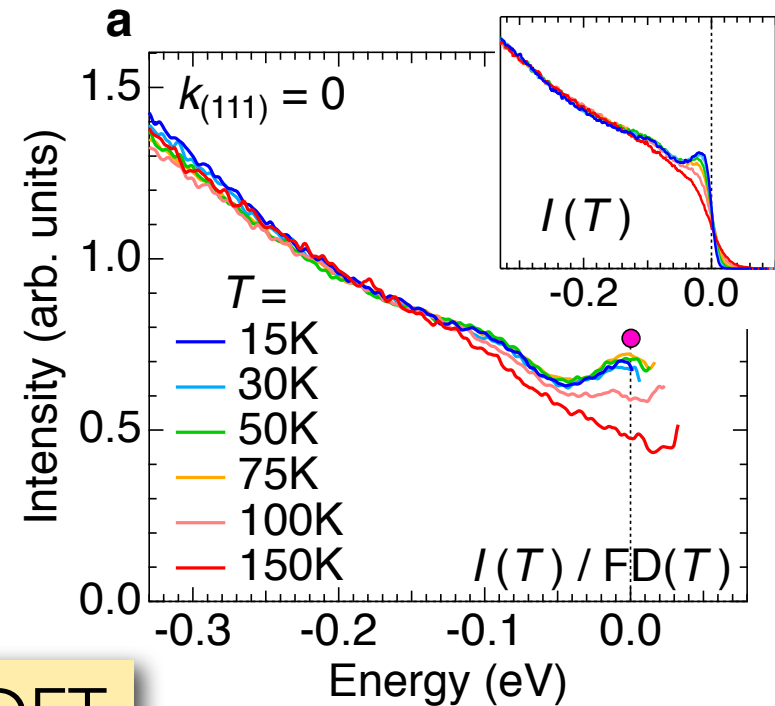
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$\text{Pr}_2\text{Ir}_2\text{O}_7$ S. Nakatsuji

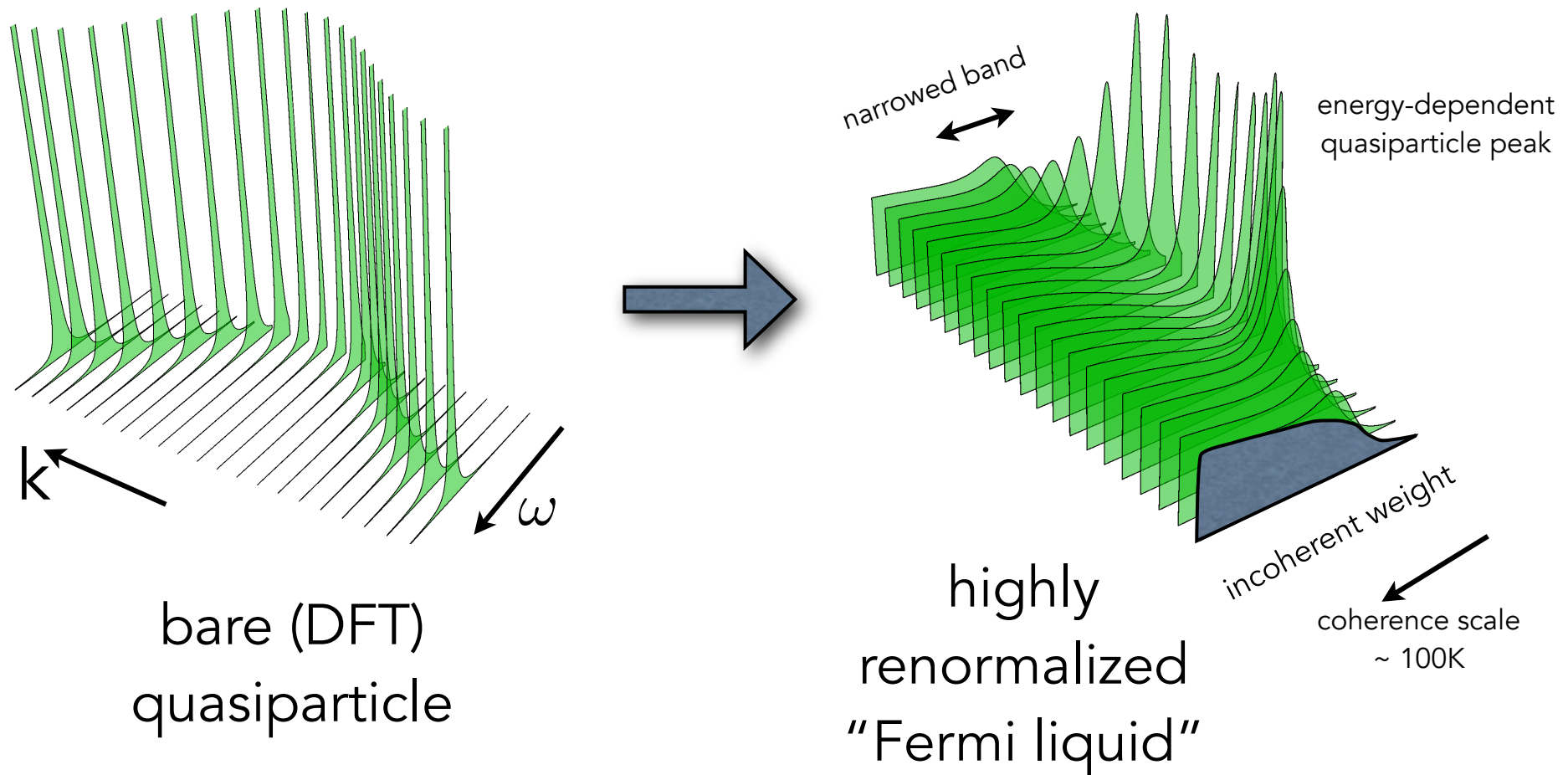
T. Kondo

S. Shin

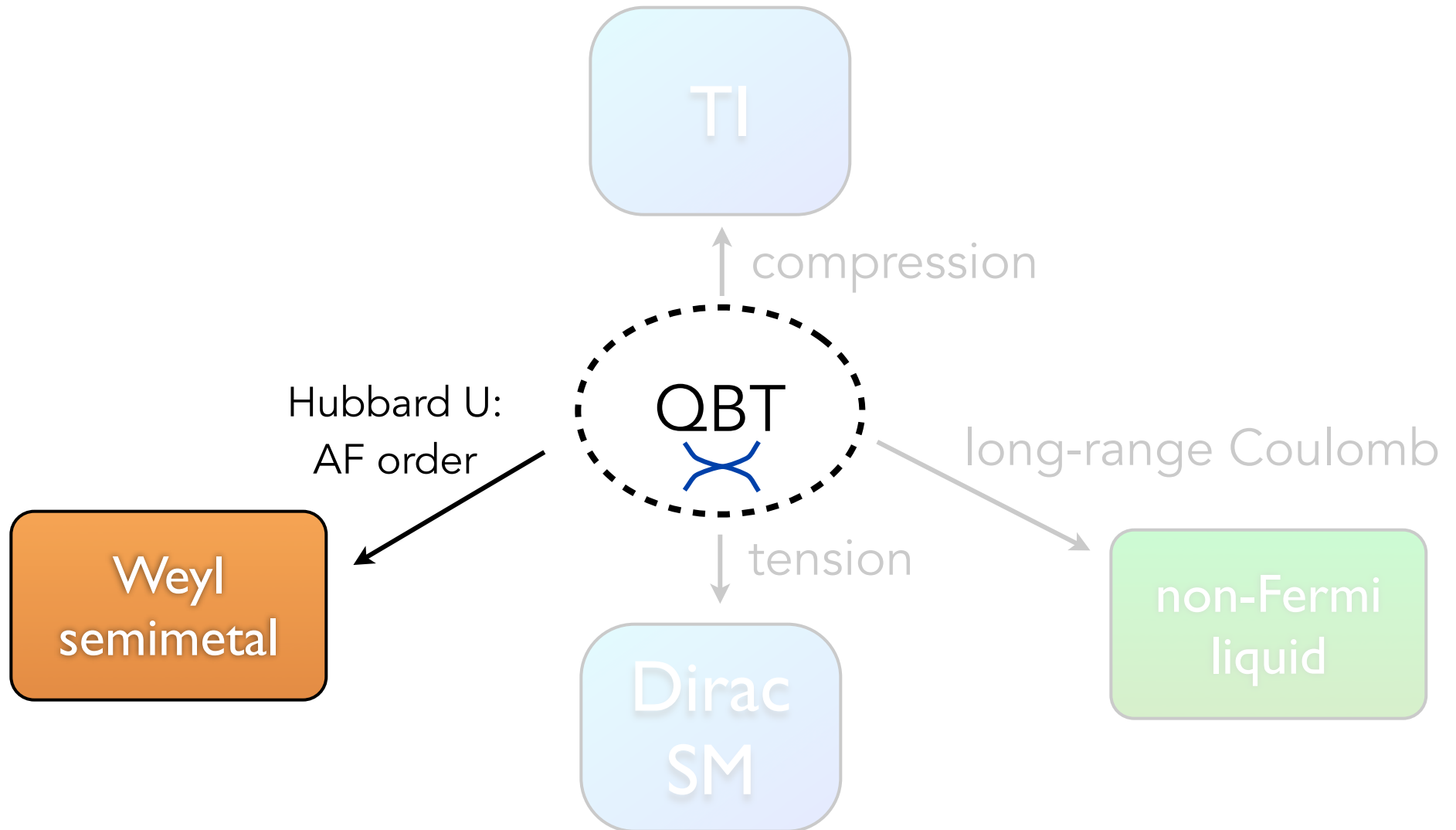


- ✓ Bandwidth reduced by 3-5 from DFT
- ✓ Quasiparticle weight strongly suppressed with temperature

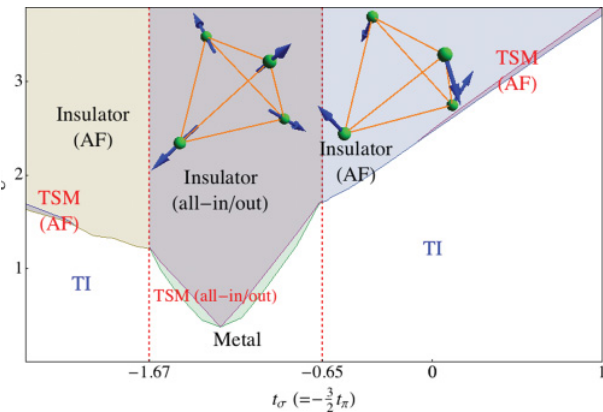
Correlated quasiparticles



Proximate phases

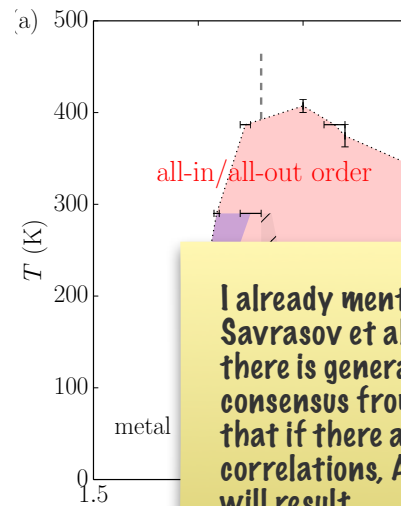


Ordering: theory



W. Witzak-Krempa + YB Kim, 2012

Hartree-Fock

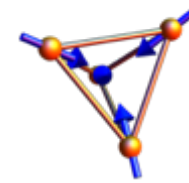
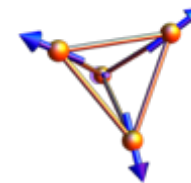


H. Sh

DMFT

Jay-Z model

$$H = J_z \sum_{\langle i,j \rangle} S_i^z S_j^z$$



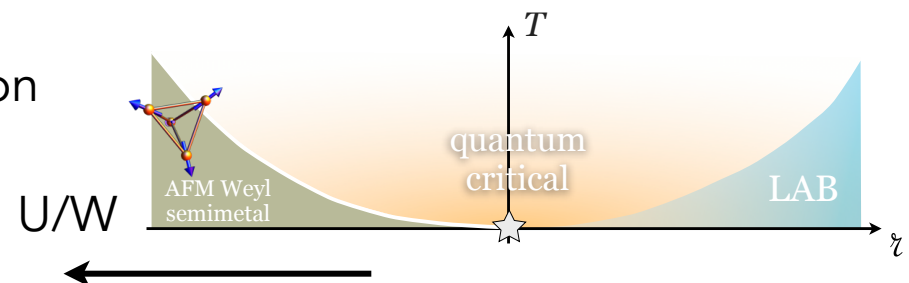
$J_z < 0$: all-in/all-out order

superexchange

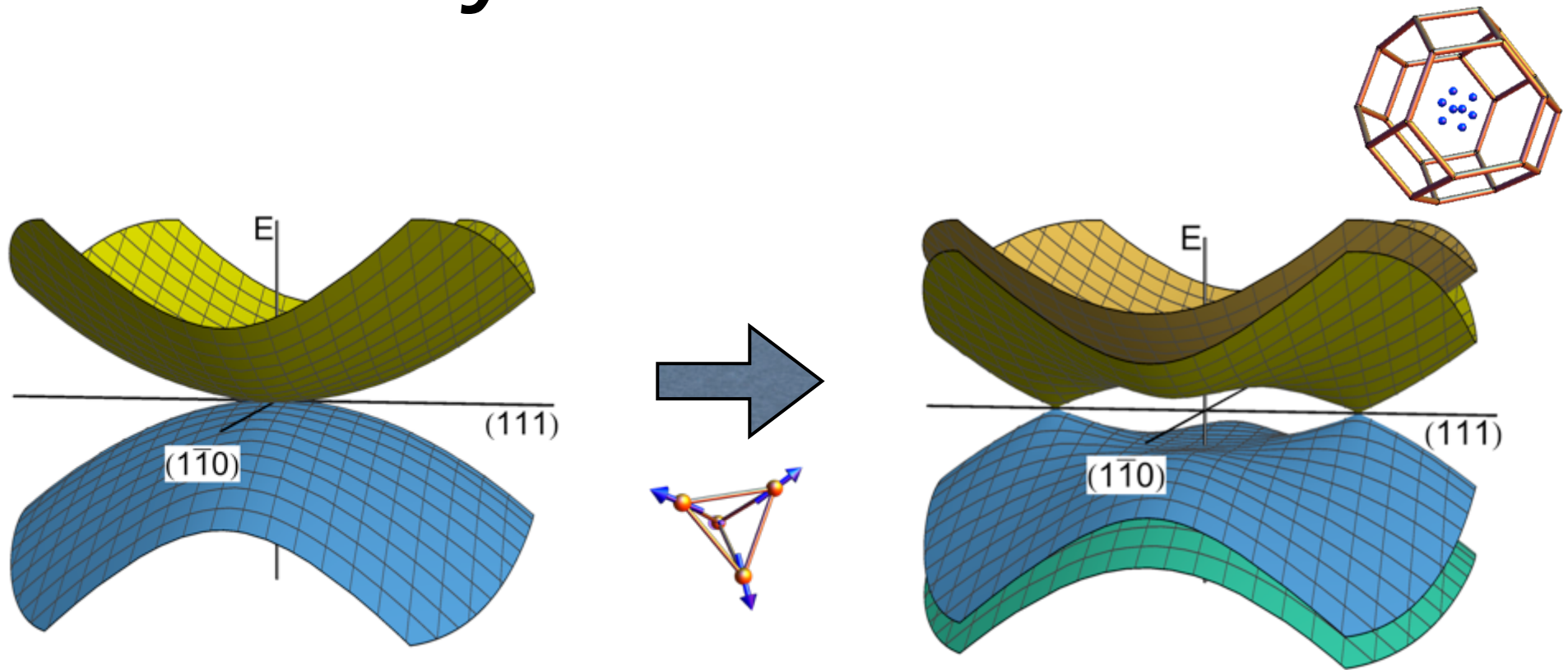
- General agreement: transition to AIAO Ising AF order

- ◆ Quantum critical theory of this transition

L. Savary et al, 2014 - topological QCP



Weyl semimetal

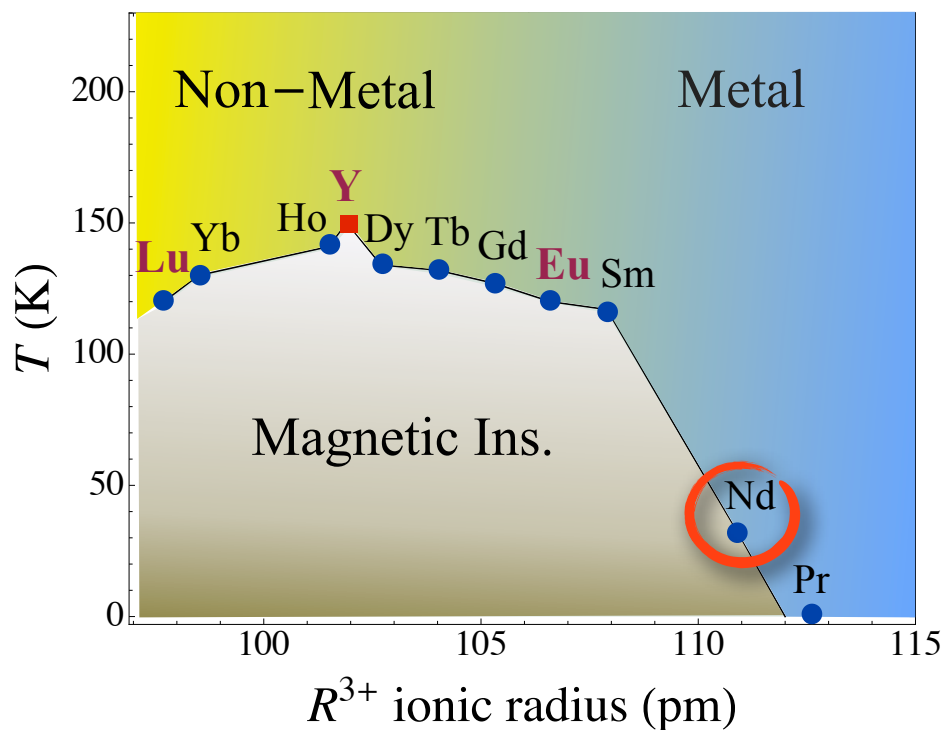


If quasiparticle picture applies at low energy, a *weakly* ordered AIAO state *must* be a Weyl semimetal

Weyl?

- Weyl semimetal?

Best candidate
should be
 $\text{Nd}_2\text{Ir}_2\text{O}_7$

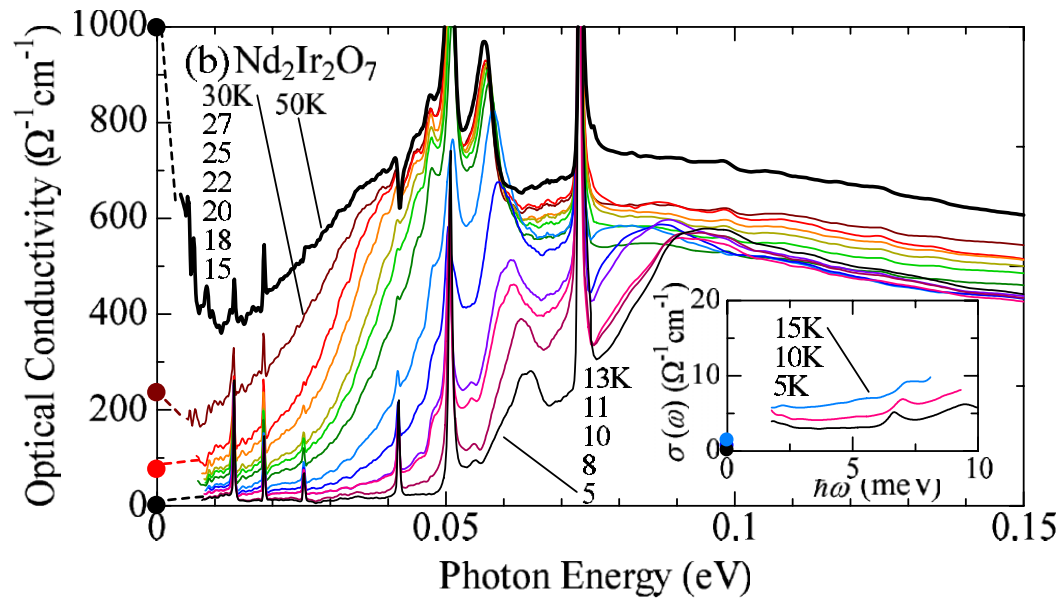


R^{3+} ionic radius (pm)
Yanagashima+Maeno, JPSJ 2001
K. Matsuhira et al, JPSJ 2011
W. Witczak-Krempa et al, ARCOMP 2013



Weyl not?

K. Ueda *et al*, 2012

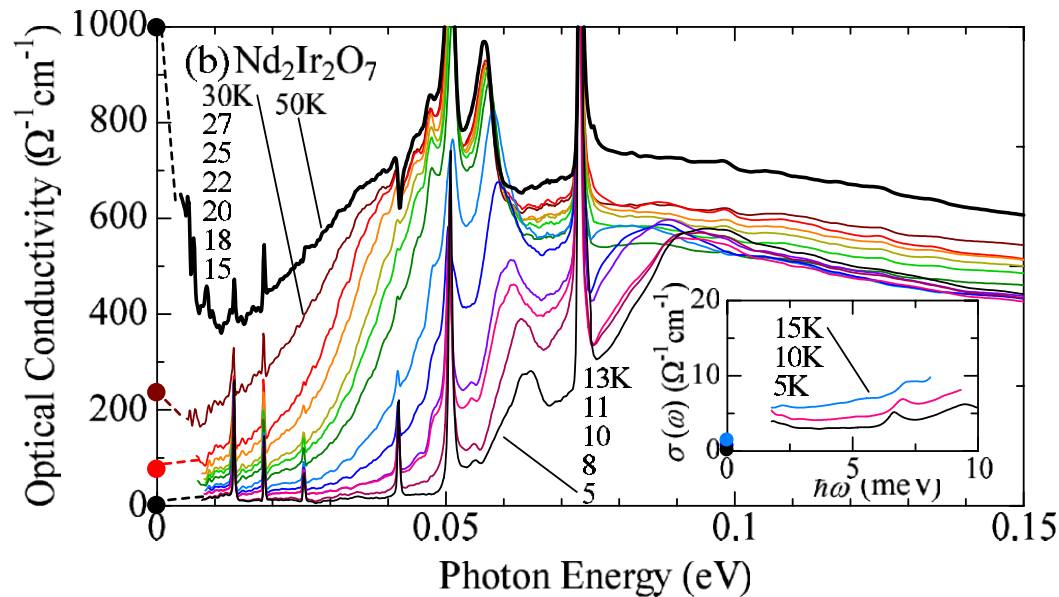


charge gap $\sim$
45meV



Weyl not?

K. Ueda *et al*, 2012



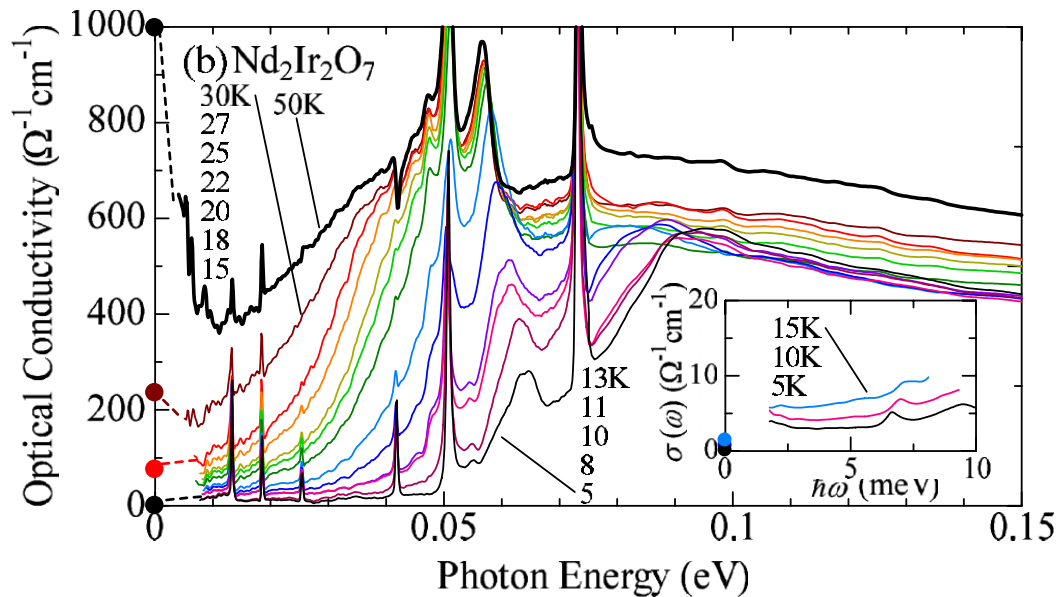
charge gap $\sim$
45meV

- ✓ Bandwidth reduced by 3-5 from DFT
- ✓ Quasiparticle weight strongly suppressed with temperature
- ✓ Gap $\sim$ 45meV $\sim$ 18 T_c



Weyl not?

K. Ueda *et al*, 2012

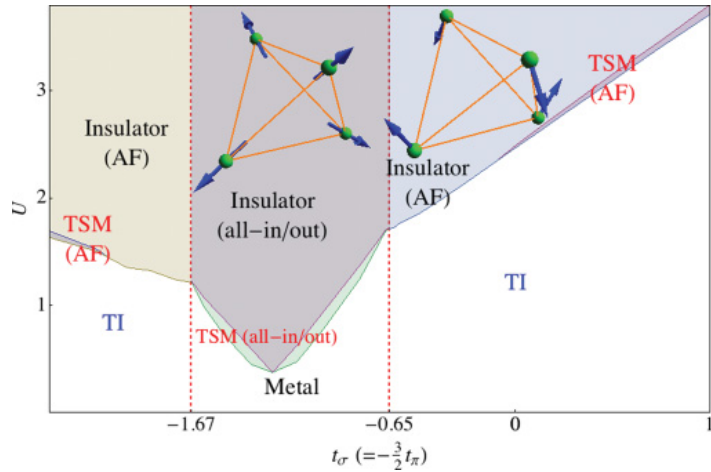


charge gap $\sim$
45 meV

If quasiparticle picture applies at low energy, a *weakly* ordered AIAO state *must* be a Weyl semimetal

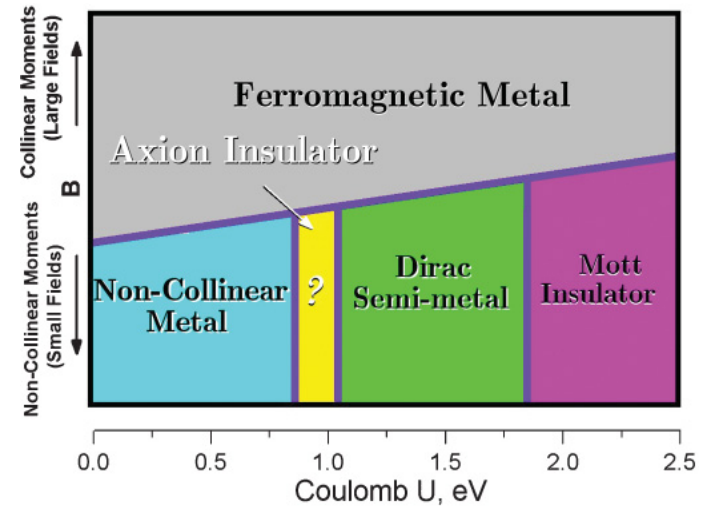
- ◆ Not weakly ordered?
- ◆ No quasiparticles?
- ◆ Weyl with very small DOS?

Weak or not?



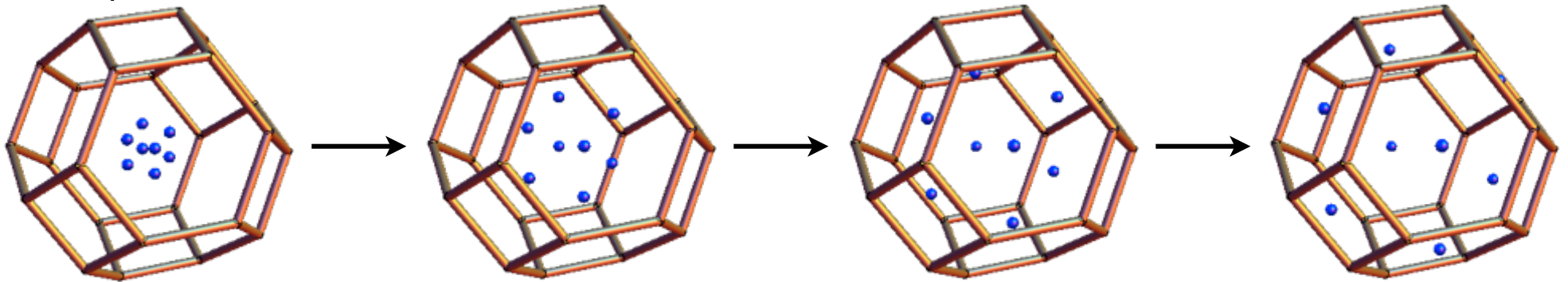
W. Witzak-Krempa + YB Kim, 2012

VS



X. Wan et al, 2011

Weyl points move to zone boundary and annihilate with increasing order?

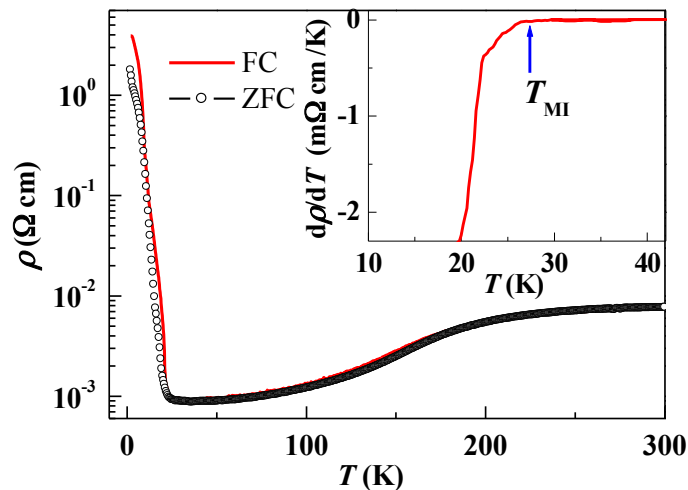


Could this have happened already for Nd?

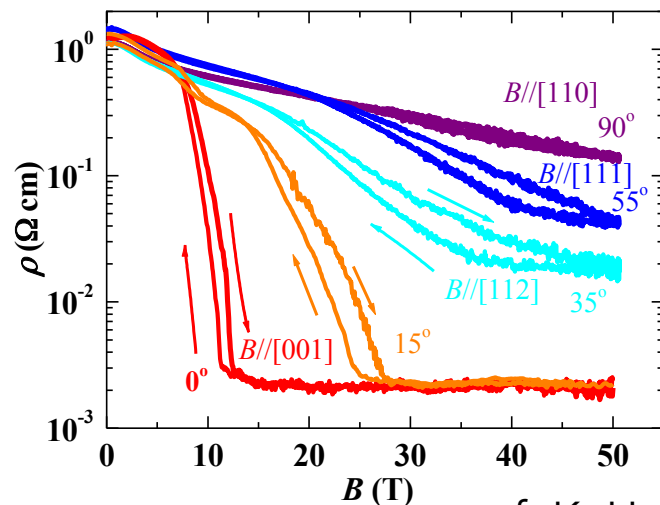


Transport

To explore this in more detail, let's look at transport



- Seems to be a transition with hard gap
- thermal transition is continuous



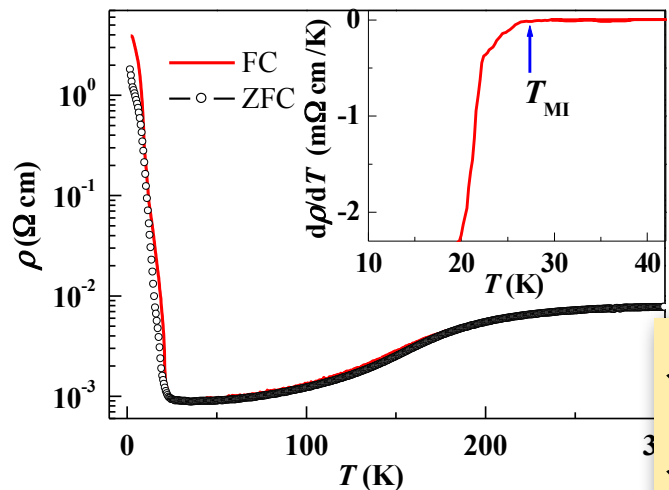
- However, abrupt first order transition appears at low T in a field
- Insulator-metal transition *only* occurs for fields along (100)

c.f. K. Ueda et al, 2015

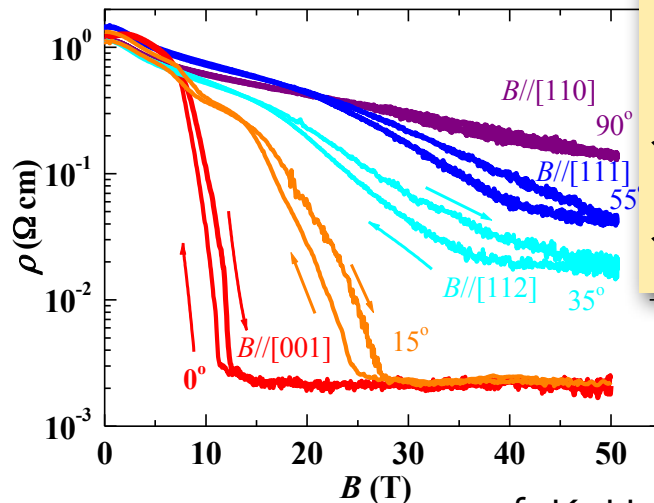


Transport

To explore this in more detail, let's look at transport



- Seems to be ... with hard gap
- thermal transport ... continuous

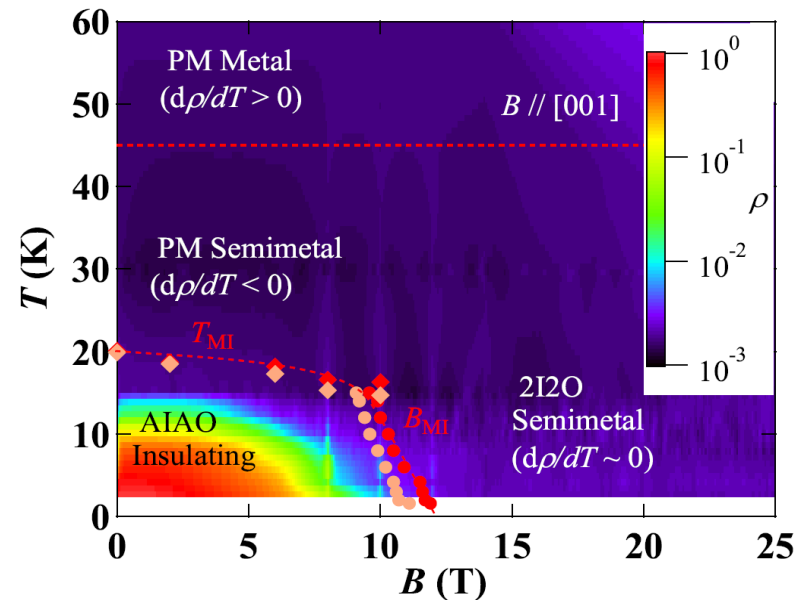
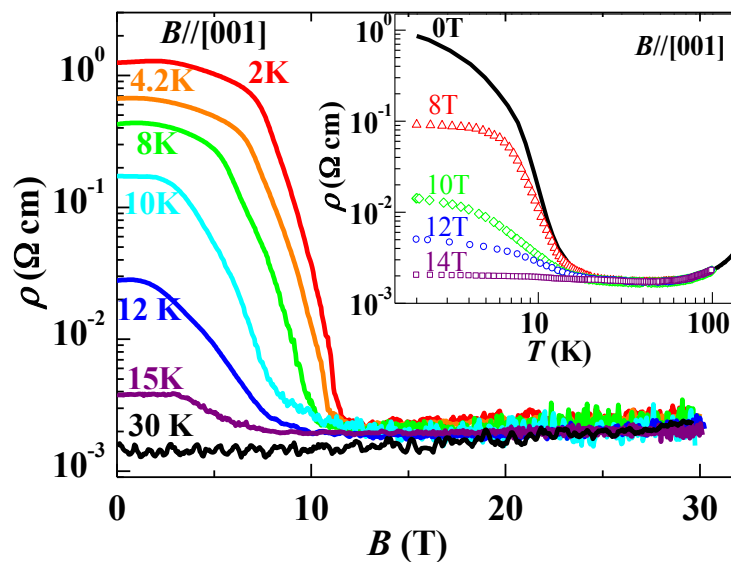


- ✓ Bandwidth reduced by 3-5 from DFT
- ✓ Quasiparticle weight strongly suppressed with temperature
- ✓ Gap $\sim 45\text{meV} \sim 18 T_c$
- ✓ MIT driven by (tiny) 12T field

fields along (100)

c.f. K. Ueda et al, 2015

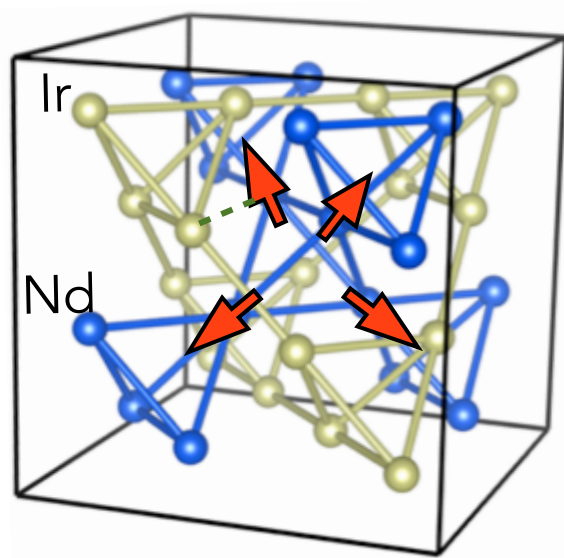
Metal-Insulator Transition



- First order transition at low temperature: mechanism of MIT is vanishing of condensation energy, not gap closing.
- Anisotropy: clear indication that Nd plays a direct role in the MIT - probably this is a unique feature for $R=\text{Nd}$

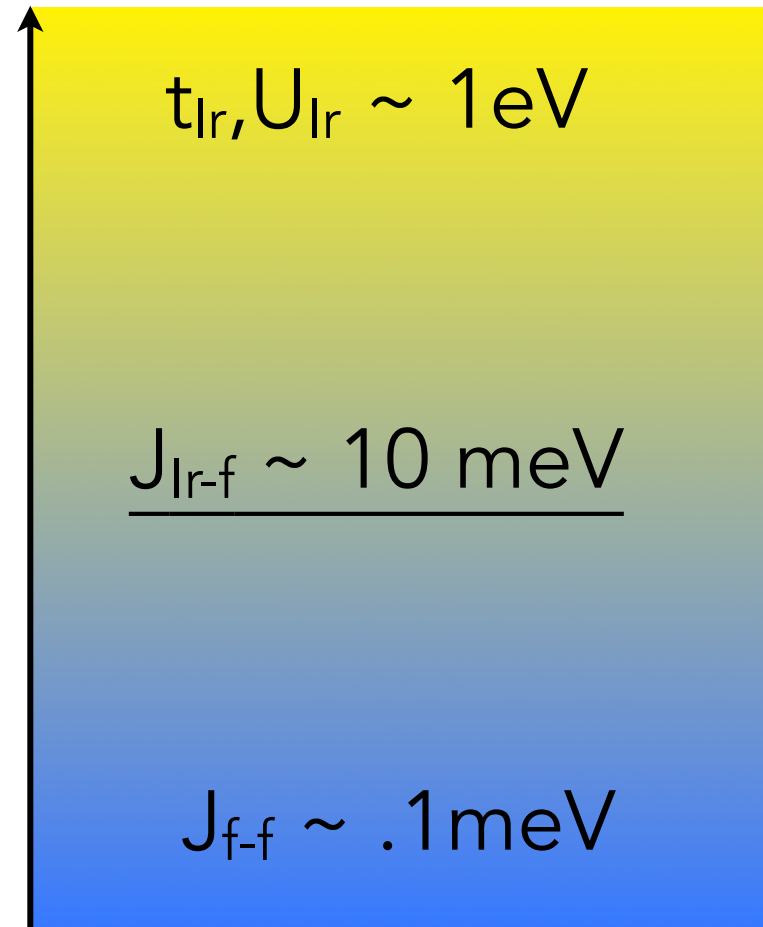
Nd physics

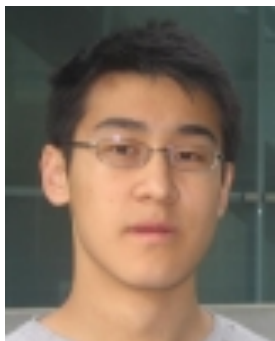
$$H_K = \sum_{ia} \vec{S}_i \cdot \mathbf{J}_{ia} \cdot \vec{\tau}_a$$



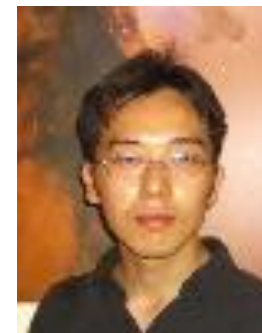
- Much of gap (45meV) is *due* to Nd-Ir exchange
- Intrinsic g-tensor anisotropy of Nd explains anisotropy of MIT

(100) field achieves maximum Nd polarization of 2in-2out type, "opposite" to AIAO Ising order



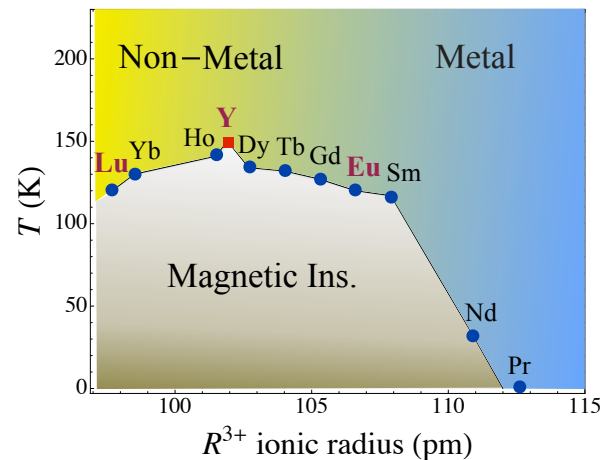
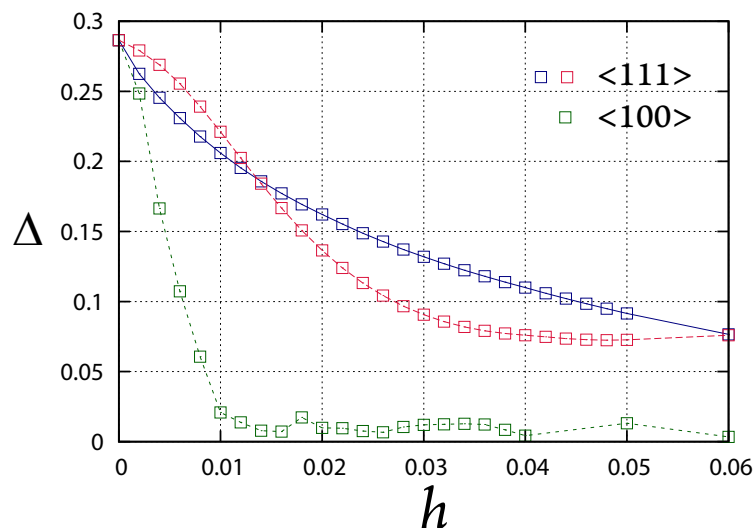


Anisotropy



- Nd sensitivity to direction is transferred to Ir

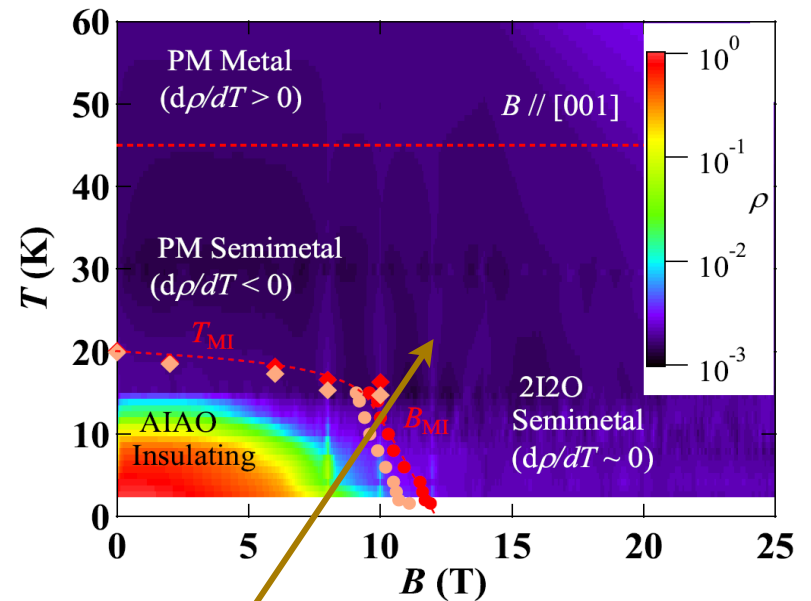
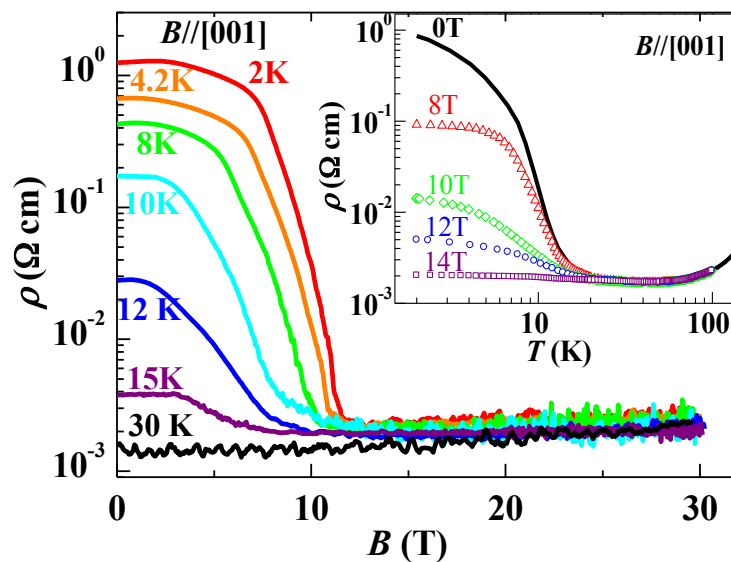
$$H_K = \sum_{ia} \vec{S}_i \cdot \mathbf{J}_{ia} \cdot \vec{\tau}_a \quad \text{acts like } h_{\text{eff}} \sim 10 \text{ meV local field on Ir}$$



Hartree-Fock on Hubbard
+Kondo lattice model

- h_{eff} is significant for Nd case
- can explain the anisotropy

Metal-Insulator Transition

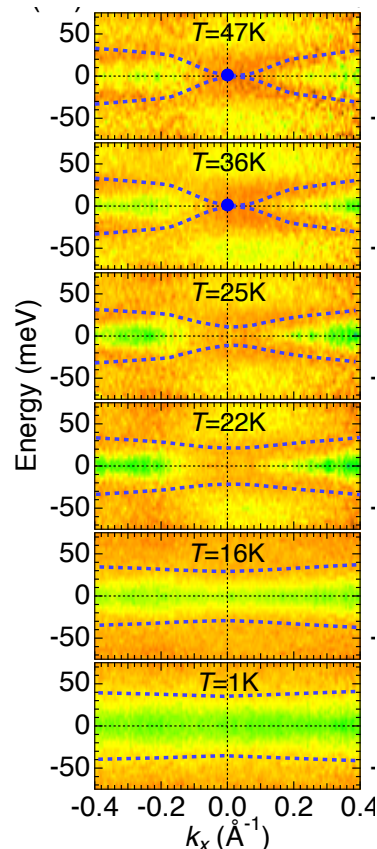


Nd moments are 2in-2out here, but no evidence that Ir is 2in-2out, or even has substantial local moments here



Quasiparticles?

- Already observed bandwidth reduction in $\text{Pr}_2\text{Ir}_2\text{O}_7$ - what happens for Nd?



$T > T_c$: like Pr

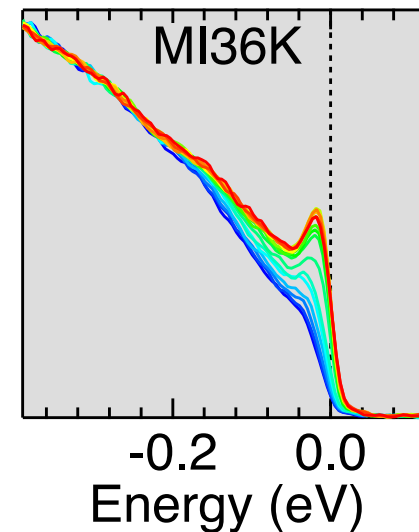
$T = T_c$: no precursor - Slater

$T < T_c$: gap developing

Slater to Mott
crossover?

$T \ll T_c$: remarkably flat

In strong Mott localized regime non-trivial band topology is unlikely



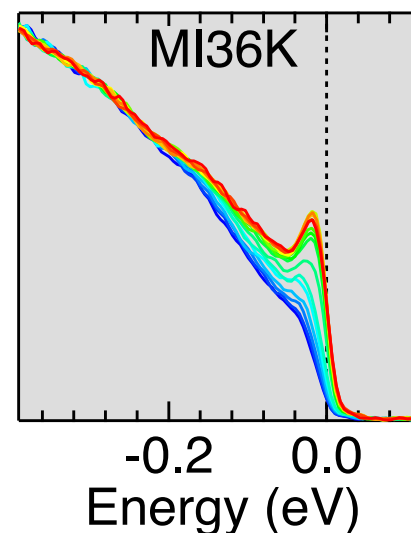
loss of quasiparticle peak



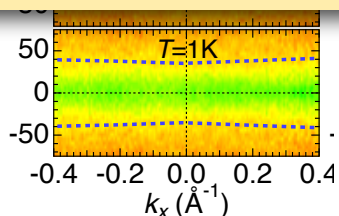
Quasiparticles?

- Already observed bandwidth reduction in $\text{Pr}_2\text{Ir}_2\text{O}_7$ - what happens for Nd?

- ✓ Bandwidth reduced by 3-5 from DFT
- ✓ Quasiparticle weight strongly suppressed with temperature
- ✓ Gap $\sim 45\text{meV} \sim 18 T_c$
- ✓ MIT driven by (tiny) 12T field
- ✓ Extreme band flattening and reduction of QP peak at low T



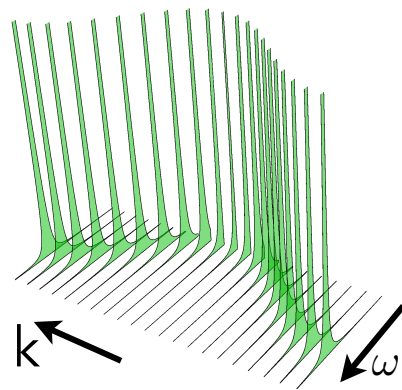
loss of quasiparticle peak



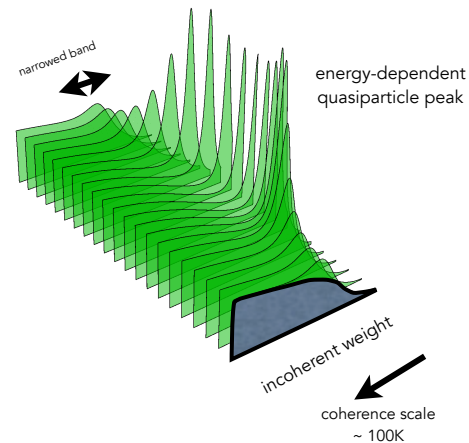
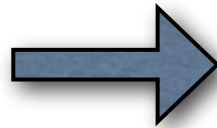
$T \ll T_c$: remarkably flat

In strong Mott localized regime non-trivial band topology is unlikely

Energy regimes



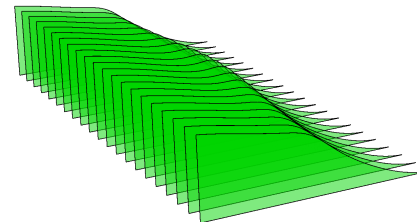
bare (DFT)
quasiparticle



highly
renormalized
"Fermi liquid"



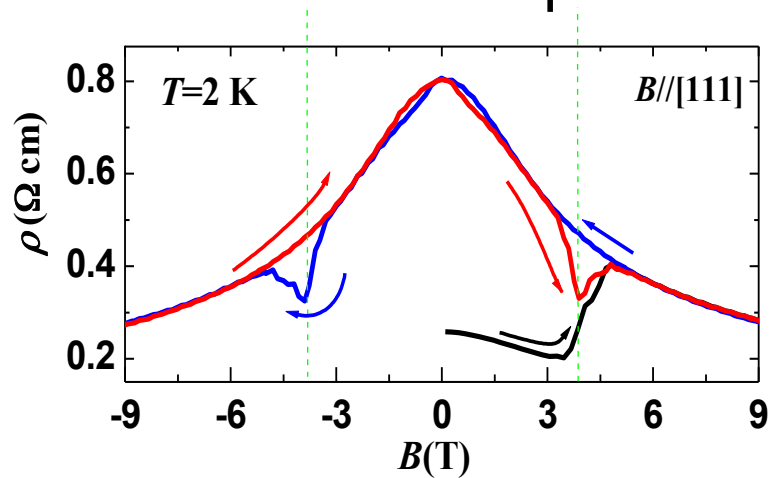
mean-field onset



Mott-like low T limit?

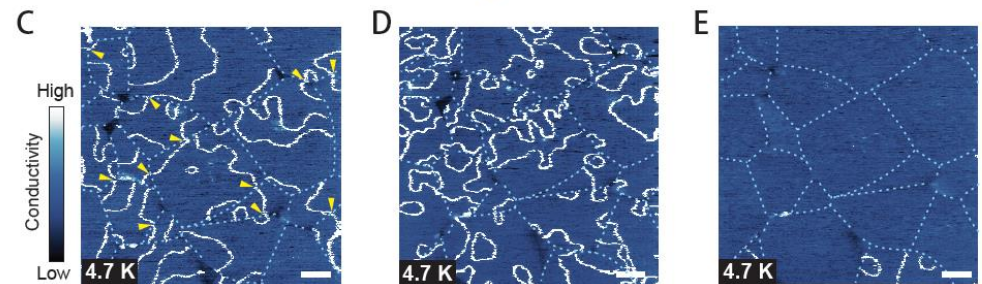
Domains

- Large magnetic moment of Nd allows manipulation of domains



Z. Tian *et al*, 2015

c.f. K. Ueda *et al* (2014)



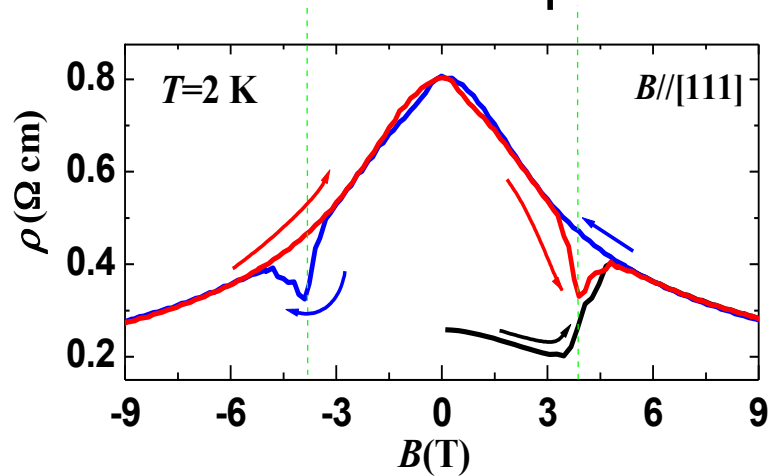
local conductivity

E. Y. Ma *et al*, 2015

conducting domain walls! c.f. previous talks

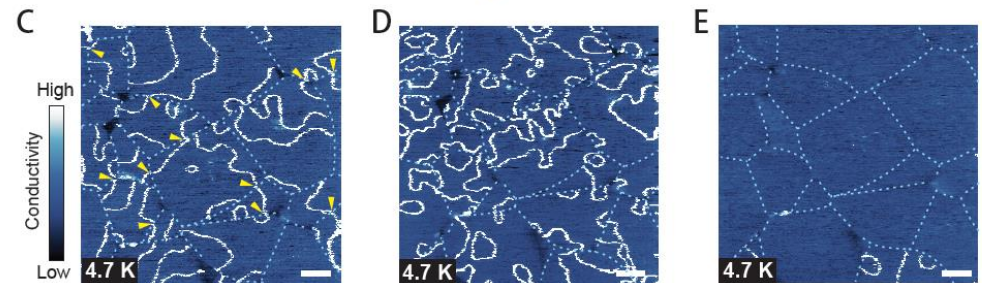
Domains

- Large magnetic moment of Nd allows manipulation of domains



Z. Tian *et al*, 2015

c.f. K. Ueda *et al* (2014)



local conductivity

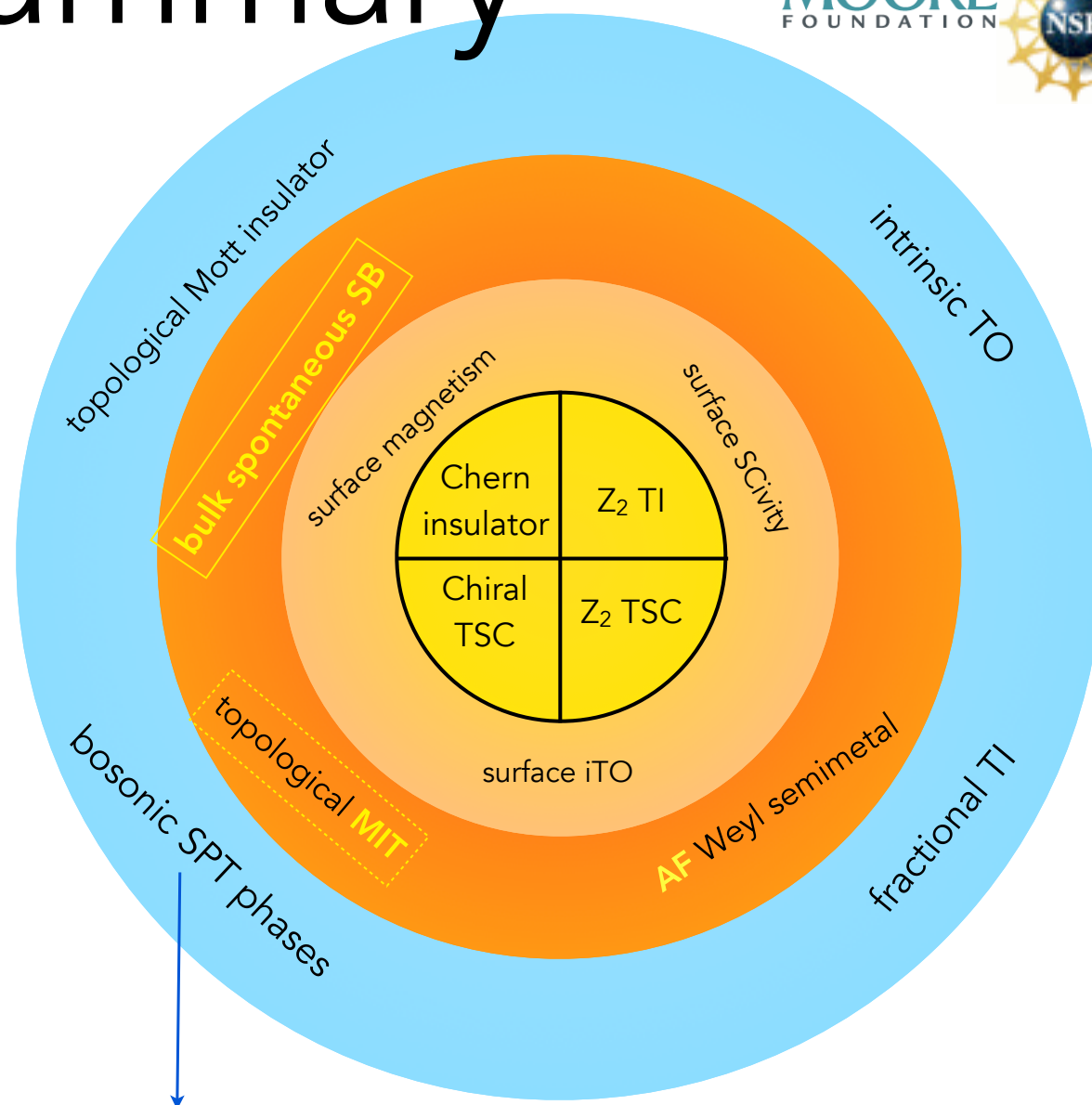
E. Y. Ma *et al*, 2015

may be evidence for correctness of Slater picture?

estimate $d \sim \rho_{\text{metal},3d} / \rho_{\square} \sim (1 \text{ m}\Omega \cdot \text{cm}) / (1 \text{ k}\Omega) \sim 10 \text{ nm}$

Summary

- Pyrochlore iridates show a rich phenomenology with nodal and gapped states, metal-insulator transitions, and strong anisotropic response
- Correlations induce antiferromagnetism but also Mott localization, which, helped by Nd f-moments, seems to subdue the Weyl fermions expected from the former alone. Challenge: can we recover them? (c.f. K. Ueda talk?)
- Quite a bit still to explore! Correlated materials prove surprising and challenging



Advertisement: Zhen Bi *et al* , arXiv:1602.03190